

## Disclaimer

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<https://www.physics-and-stuff.com/>

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31.05.16

Condensed Matter Exercise 3

2.a. Quantum Well: 
$$V(x) = \begin{cases} V_0 & \text{for } -L/2 \leq x \leq L/2 \\ \infty & \text{else} \end{cases}$$

S.E yields: 
$$-\frac{\hbar^2}{2m} \nabla^2 \psi(\vec{r}) + V(x) \psi(\vec{r}) = E \psi(\vec{r})$$

with Ansatz: 
$$\psi(\vec{r}) = \phi_n(x) e^{i(k_y y + k_z z)} = \phi_n(x) e^{i\vec{k}_n \cdot \vec{r}}$$

$$\Rightarrow -\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) \phi_n(x) e^{i\vec{k}_n \cdot \vec{r}} + V(x) \phi_n(x) e^{i\vec{k}_n \cdot \vec{r}} = E \phi_n(x) e^{i\vec{k}_n \cdot \vec{r}}$$

Separation of variables

$$\Rightarrow -\frac{\hbar^2}{2m} \left( \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) e^{i\vec{k}_n \cdot \vec{r}} = E_n e^{i\vec{k}_n \cdot \vec{r}} \quad (1)$$

and 
$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \phi_n(x) + V(x) \phi_n(x) = E_x \phi_n(x) \quad (2)$$

(1) 
$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial y^2} e^{i k_y y} = E_y e^{i k_y y} \quad \Bigg| \quad -\frac{\hbar^2}{2m} \frac{\partial^2}{\partial z^2} e^{i k_z z} = E_z e^{i k_z z}$$

$$\Leftrightarrow \frac{\hbar^2}{2m} k_y^2 = E_y \quad \Bigg| \quad \Leftrightarrow \frac{\hbar^2}{2m} k_z^2 = E_z$$

$$\Rightarrow E_n = \frac{\hbar^2}{2m} \vec{k}_n^2$$

(2) 
$$-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial x^2} \phi_n(x) = (E_x - V(x)) \phi_n(x)$$

$$\phi_n(x) = 0 \quad \text{for } x < -L/2, x > L/2$$

$$V(x) = V_0 \quad \text{for } -L/2 \leq x \leq L/2$$

$$\Rightarrow \frac{\partial^2}{\partial x^2} \phi_n(x) = -\frac{2m(E_x - V_0)}{\hbar^2} \phi_n(x)$$

Ansatz: 
$$\phi_n(x) = e^{\lambda x} \Rightarrow \phi_n(x) = x^2 e^{\lambda x}$$

$$\lambda^2 = -\frac{2m(E_x - V_0)}{\hbar^2} \Leftrightarrow \lambda = \pm i \sqrt{\frac{2m(E_x - V_0)}{\hbar^2}} = \pm i k_x$$

$$\Rightarrow \phi_n(x) = A e^{i k_x x} + B e^{-i k_x x}$$

boundary conditions: 
$$\phi_n(-L/2) = \phi_n(L/2) = 0$$

$$\phi_n(-L/2) = A e^{-i k_x L/2} + B e^{i k_x L/2} = 0$$

$$\Leftrightarrow A = -B e^{i k_x L}$$

Ansatz;  
Sep. of variables!  
 $\phi_1 \phi_2 \phi_3 \rightarrow A+B+C = E_{\text{total}}$   
 $\vec{k}_y$ ?  
parallel to  
the potential  
barriers  
 $V(y) = V(z) = 0$   
constant potential  
in  $y$  and  $z$   
 $\rightarrow$  shell  
Ansatz for  
this equation?  
No wavefunction  
split in product?  
 $\leftarrow$  sep. is split  
in product!



$$\psi_n(x/2) = A e^{ikx/2} + B e^{-ikx/2} = 0$$

$$\Leftrightarrow A + B e^{-ikxL} = 0$$

$$\Leftrightarrow B (e^{-ikxL} - e^{ikxL}) = 0$$

$$\Leftrightarrow -B \cdot 2i \sin(kxL) = 0$$

$$\Leftrightarrow \sin(kxL) = 0 \Leftrightarrow kxL = n\pi \Leftrightarrow k_x = \frac{n\pi}{L} = \frac{\sqrt{2m(E_x - V_0)}}{\hbar}$$

$$\Leftrightarrow E_x' = \frac{\hbar^2 n^2 \pi^2}{2mL^2} = E_x - V_0$$

$$\Rightarrow \text{Total Energy yields: } E(\vec{k}_n, n) = \frac{\hbar^2}{2m} k_n^2 + \frac{\hbar^2 n^2 \pi^2}{2mL^2}$$

$$\text{Volume of } k\text{-space: } \Delta_k = \frac{\pi^2}{L^2}$$

2D Sub band: minimum  $E(0, n)$  for  $\vec{k}_n = 0$

$$dN = \frac{2}{4} 2\pi k \frac{dk}{\frac{\pi^2}{L^2}} = \frac{L^2}{\pi} k dk, \quad k (= \frac{n\pi}{L}) = \frac{\sqrt{2m(E')}}{\hbar}$$

$$= \frac{L^2}{\pi} \frac{\sqrt{2mE'}}{\hbar} \frac{dk}{dE} dE = \frac{L^2}{\pi} \frac{\sqrt{2mE'}}{\hbar} \frac{1}{2} \frac{1}{\sqrt{2mE'}} \frac{1}{\hbar} dE$$

$$= \frac{L^2}{\pi} \frac{1}{\hbar^2} \frac{1}{2} \sqrt{4m^2 E'} \cdot \frac{1}{\hbar} dE = \frac{L^2 m}{\hbar^2 \pi} dE$$

$$\frac{dN}{dE} = \frac{m}{\hbar^2 \pi} dE \Rightarrow D(E) = \frac{m}{\hbar^2 \pi}$$



Free + bound states?  
bound in 1-dimension  
free in 2-dim

Why  $E_x'$ ?  
not  $E_k'$ ?  
total energy

Subband minimum?  
fixed in (new plane with energy levels) and first energy level of states

$\Delta_k = \frac{\pi^2}{L^2}$ ?  
k-volume

$(L^2) = V = L_x L_y$   
Volume  
E of confined area in 2D



b) Quantum Wire 1D

$$V(x,y) = \begin{cases} V_0 & \text{for } -\frac{L}{2} \leq x \leq \frac{L}{2} \text{ and } -\frac{L}{2} \leq y \leq \frac{L}{2} \\ 0 & \text{else} \end{cases}$$

Calculation analogous to a) but  $\Delta k = \frac{\pi}{L}$  (k-space volume)

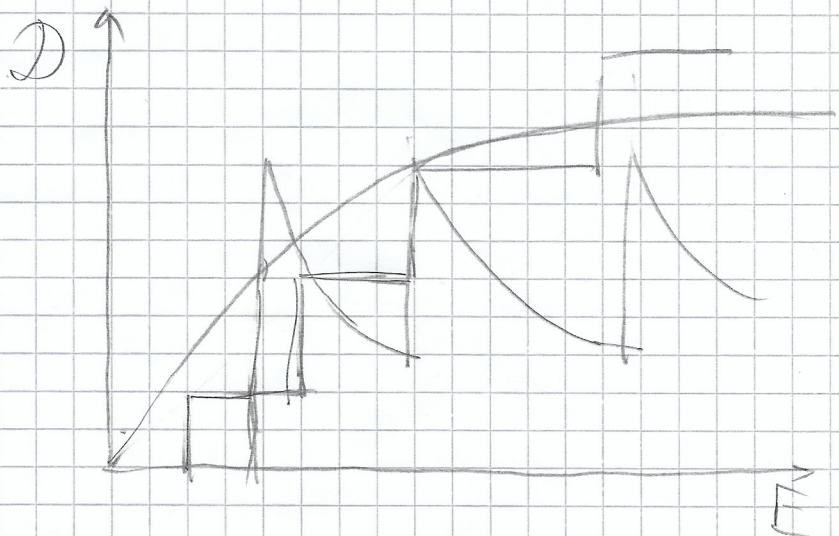
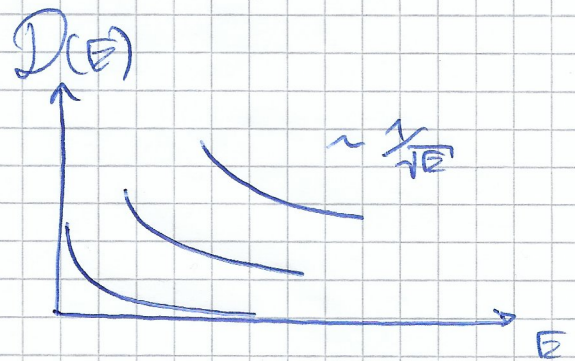
$$E = \frac{\hbar^2 \alpha^2}{2mL^2} (n_x^2 + n_y^2) + \frac{\hbar^2 k_z^2}{2m} \quad L \rightarrow L_y \text{ and } L_z$$

2 for spin  
but  $\frac{1}{2}$   
positive and negative  
numbers have  
to be  
taken  
out of  
calculus

$$\begin{aligned} dN &= \frac{2}{2} \frac{dk}{\pi/L} = \frac{L}{\pi} dk \quad \text{with } k = \sqrt{\frac{2mE}{\hbar^2}} \\ &= \frac{L}{\pi} \frac{dk}{dE} dE = \frac{L}{\pi} \frac{1}{\hbar} \sqrt{m} \cdot \frac{1}{2} \frac{1}{\sqrt{E}} dE \\ &= \frac{L}{\pi} \frac{1}{\hbar} \sqrt{\frac{m}{2E}} dE \end{aligned}$$

$$\Rightarrow \frac{dN}{V} = \frac{1}{\pi \hbar} \sqrt{\frac{m}{2E}} dE$$

$$\Rightarrow D(E) = \frac{1}{\pi \hbar} \sqrt{\frac{m}{2E}}$$





#### 4. Chemical potential in lower dimensions

$$n = \int_0^{\infty} D(E) f(E) dE \quad \text{with} \quad D(E) = \frac{m}{\pi \hbar^2} \quad \text{in 2D}$$

Where from  
 $D(E), f(E)$ !

$$f(E) = \frac{1}{e^{\frac{E-\mu}{k_B T}} + 1} \quad \text{Fermi-Distribution}$$

$$n = \frac{m}{\pi \hbar^2} \int_0^{\infty} \frac{1}{e^{\frac{E-\mu}{k_B T}} + 1} dE = \frac{m}{\pi \hbar^2} \int_0^{\infty} \frac{e^{-\frac{(E-\mu)}{k_B T}}}{1 + e^{-\frac{(E-\mu)}{k_B T}}} dE$$

$$= - \frac{m k_B T}{\pi \hbar^2} \left. \ln(1 + e^{-\frac{(E-\mu)}{k_B T}}) \right|_0^{\infty}$$

$$= + \frac{m k_B T}{\pi \hbar^2} \ln(1 + e^{\frac{\mu}{k_B T}})$$

Solving for  $\mu$  yields:

$$\frac{n \pi \hbar^2}{m k_B T} = \ln(1 + e^{\frac{\mu}{k_B T}})$$

$$\Leftrightarrow e^{\frac{\mu}{k_B T}} = e^{\frac{k_B T n \pi \hbar^2}{m k_B T}} - 1$$

$$\Leftrightarrow \mu = k_B T \ln\left(e^{\frac{k_B T n \pi \hbar^2}{m k_B T}} - 1\right), \quad \hbar = \frac{h}{2\pi}$$

#### 5. Fermi-Energy and Fermi-Temperatur

$$E_F = \frac{\hbar^2}{2m} (3 \pi^2 n)^{2/3}, \quad T_F = \frac{E_F}{k_B}$$

$$a) \rho = \frac{m}{V} = \frac{N \cdot m_{Au}}{V} \Leftrightarrow n = \frac{N}{V} = \frac{\rho}{m_{Au}}$$

$$n = \frac{19 \frac{\text{g}}{\text{cm}^3}}{197 \text{u}} = \frac{19 \cdot 10^3 \frac{\text{kg}}{\text{cm}^3}}{3,271 \cdot 10^{-25} \text{kg}} = 5,81 \cdot 10^{22} \frac{1}{\text{cm}^3}$$

$$= 5,81 \cdot 10^{28} \frac{1}{\text{m}^3}$$

$$\Rightarrow E_F =$$

$$8,764 \cdot 10^{-19} \text{ J} = 5,5 \text{ eV}$$

$$T_F =$$

$$63507,246 \text{ K}$$



$$5) \quad n = \frac{\rho}{m_{\text{HCl}_3}} = \frac{0,081 \frac{\text{g}}{\text{cm}^3}}{34} = 1,62 \cdot 10^{-22} \frac{1}{\text{cm}^3} = 1,62 \cdot 10^{22} \frac{1}{\text{m}^3}$$

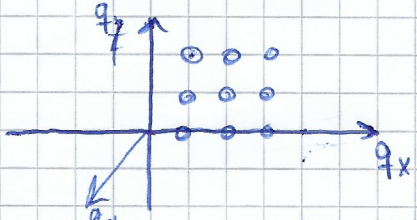
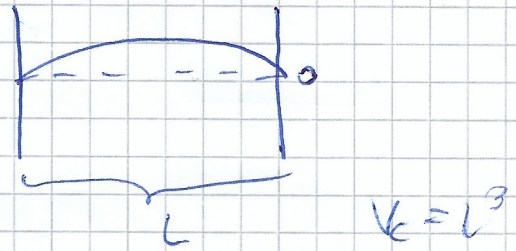
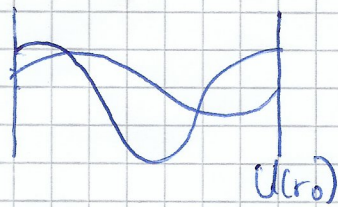
$$\Rightarrow E_F = 4,2 \cdot 10^{-4} \text{ eV}$$

$$T_F = 4,88 \text{ K}$$

$\rho_{\text{HCl}_3}$ ?



# 1. Density of States of Phonons



$$q_i = \frac{\pi}{L}, \frac{2\pi}{L}, \frac{3\pi}{L}, \dots$$

$$D(\omega) = \frac{Vc}{\pi^3} \int_{S^+(\omega)} \frac{dS_{\omega}}{|\nabla_{\mathbf{q}} \omega(\mathbf{q})|} = \frac{1}{8} \frac{Vc}{\pi^3} \int_{S(\omega)} \frac{dS_{\omega}}{|\nabla_{\mathbf{q}} \omega(\mathbf{q})|}$$

only positive

$$\frac{Vc}{(2\pi)^3} \int_{S(\omega)} \frac{dS_{\omega}}{|\nabla_{\mathbf{q}} \omega(\mathbf{q})|}$$

Same result as for periodic boundary conditions

# 3. Fermi-Dirac Distribution

$$f(E) = \frac{1}{\exp[(E-\mu)/k_B T] + 1}$$

$$f(E)|_{\mu} = f(\mu) + \left. \frac{\partial f}{\partial E} \right|_{E=\mu} (E-\mu)$$

$$= \frac{1}{2} - \frac{\exp((E-\mu)/k_B T)}{k_B T (\exp(\mu/k_B T) + \exp(E/k_B T))^2} (E-\mu)$$

$$= \frac{1}{2} - \frac{\exp(2\mu/k_B T)}{k_B T \cdot 4 \exp(2\mu/k_B T)} (E-\mu) = \frac{1}{2} - \frac{E-\mu}{4k_B T}$$

$$f(E_1) = 1, f(E_2) = 0$$

$$\frac{1}{2} = \frac{E_1 - \mu}{4k_B T} \quad \Delta E = 4k_B T$$

$$E_2 = 2k_B T + \mu$$

$$\frac{1}{2} = \frac{E_1 - \mu}{4k_B T} \quad \frac{\Delta E}{E - \mu} \approx \frac{\Delta E}{\mu} = \frac{4k_B T}{\mu}$$

$$-2k_B T + \mu = E_1$$

