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Condensed Matter Exercise 5

28.06.16

1. Electrons in homogeneous Magnetic Fields
 $\vec{E}=0$, hence the eq. of motion yields with $\vec{F}=\vec{F}_L$

$$m\ddot{\vec{r}} = -e(\dot{\vec{r}} \times \vec{B}) = m\ddot{\vec{r}} = -e(\dot{\vec{r}} \times \vec{B})$$

Rotating our coordinates, so that $\vec{B} = B\vec{e}_z$ and

$$\vec{r} = \begin{pmatrix} x \\ y \\ z \end{pmatrix}, \quad \dot{\vec{r}} = \begin{pmatrix} \dot{x} \\ \dot{y} \\ \dot{z} \end{pmatrix}$$

$$\Rightarrow \dot{\vec{r}} = -\frac{e}{m}(\dot{\vec{r}} \times \vec{B}) = -\frac{eB}{m}(\dot{\vec{r}} \times \vec{e}_z)$$

$$= -\frac{eB}{m} \begin{pmatrix} \dot{y} \\ -\dot{x} \\ 0 \end{pmatrix} = +\frac{eB}{m} (\dot{x}\vec{e}_y - \dot{y}\vec{e}_x)$$

\Rightarrow results in 3 coupled diff. eq.

$$\ddot{x} = -\frac{e}{m} B \dot{y} \quad (1)$$

$$\ddot{y} = \frac{e}{m} B \dot{x} \quad (2)$$

$$\ddot{z} = 0 \quad (3)$$

CALCULATION LIKE THIS NOT VALID BECAUSE WE DO NOT CONSIDER BAND STRUCTURE

$$(3) \Rightarrow z(t) = C_1 t + C_2, \quad C_i = \text{const}, \quad i=1,2$$

$$(1) \Rightarrow \ddot{x} = -\frac{e}{m} B \dot{y} = -\frac{e^2}{m^2} B^2 x$$

$$(2) \Rightarrow \ddot{y} = \frac{e}{m} B \dot{x} = -\frac{e^2}{m^2} B^2 y$$

$$\left. \begin{aligned} &A \{ \cos(\omega t) + i \sin(\omega t) \} \\ &+ B \{ \cos(\omega t) - i \sin(\omega t) \} \end{aligned} \right\}$$

$$x(t) = e^{\lambda t} \Rightarrow \dot{x}(t) = \lambda e^{\lambda t} \quad \text{and} \quad \ddot{x}(t) = \lambda^2 e^{\lambda t}$$

$$\lambda^2 e^{\lambda t} = -\frac{e^2}{m^2} B^2 \lambda e^{\lambda t} \Rightarrow \lambda^2 = -\frac{e^2}{m^2} B^2$$

$$\Rightarrow x(t) = A e^{i\frac{eB}{m}t} + B e^{-i\frac{eB}{m}t}$$

Analogue $y(t) = A e^{i\frac{eB}{m}t} + B e^{-i\frac{eB}{m}t}$

$$z(t) = C_1 t + C_2$$

$$\begin{aligned} \Rightarrow \vec{r}_\perp &= \vec{r} - \vec{B}(\vec{B} \cdot \vec{r}) = \vec{r} - \vec{e}_z(\vec{e}_z \cdot \vec{r}) = \vec{r} - \vec{e}_z z(t) \\ &= x(t)\vec{e}_x + y(t)\vec{e}_y \end{aligned}$$

Why are these the trajectories in k-space? lecture

Bandwidth? yes

Only real parts for exp?

(linear combination of sinus + cosinus, so that we get real part)

Board:

$$\frac{dk}{dt} = -\frac{e}{\hbar} \left(\frac{dr}{dt} \times \underline{B} \right)$$

$$k(t) - k(0) = -\frac{e}{\hbar} (r(t) - r(0)) \times \underline{B}$$

$$\underline{B} = B \underline{e} = -\frac{e}{\hbar} (r(t) - r(0))_{\perp} \times \underline{B}$$

$$x(t) - x(0) = \frac{\hbar}{eB} (k_y(t) - k_y(0))$$

$$y(t) - y(0) = -\frac{\hbar}{eB} (k_x(t) - k_x(0))$$

Without assuming $B = B \underline{e}$ · $B \times$ on both sides

$$\underline{B} \times \dot{\underline{r}} = \underline{B} \times \left(-\frac{e}{\hbar} \left(\frac{dr}{dt} \times \underline{B} \right) \right) = -eB \dot{\underline{r}}_{\perp}$$

2. Classical theory of Cyclotron Resonance

isotropic effective mass? / relaxation time?

$$m^* \dot{\vec{v}} + \frac{m^*}{\tau} \vec{v} = q(\vec{E} + \vec{v} \times \vec{B})$$

Ansatz: $\vec{E} = \vec{E}_0 e^{i\omega t}$, $\vec{E}_0 = \begin{pmatrix} E_x \\ 0 \\ 0 \end{pmatrix}$

Otherwise consider Bandstructure

Ansatz? / rel. time: time between Scatters etc.
Driven by E-field with c

$$\vec{v}(t) = \vec{v}_0 e^{i\omega t} \Rightarrow \vec{v}(t) = i\omega \vec{v}_0 e^{i\omega t}, \vec{v}_0 = \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix}$$

$$\vec{B} = \begin{pmatrix} 0 \\ 0 \\ B_z \end{pmatrix}$$

yields: $i\omega \vec{v}_0 + \frac{1}{\tau} \vec{v}_0 = \frac{q}{m^*} (\vec{E}_0 e^{i\omega t} + (\vec{v}_0 \times \vec{B}) e^{i\omega t})$

$$\Leftrightarrow \vec{v}_0 \left(i\omega + \frac{1}{\tau} \right) = \frac{q}{m^*} (\vec{E}_0 + (\vec{v}_0 \times \vec{B}))$$

$$\Leftrightarrow (1 + i\omega\tau) \vec{v}_0 = \frac{q\tau}{m^*} (\vec{E}_0 + (\vec{v}_0 \times \vec{B}))$$

$$\Leftrightarrow \begin{pmatrix} v_x \\ v_y \\ v_z \end{pmatrix} = \frac{q\tau}{(1+i\omega\tau)m^*} \left[\begin{pmatrix} E_x \\ 0 \\ 0 \end{pmatrix} + \begin{pmatrix} v_y B_z \\ -v_x B_z \\ 0 \end{pmatrix} \right]$$

$$v_x = \frac{q\tau}{m^*} (E_x + v_y B_z)$$

$$v_y = -\frac{q\tau}{m^*} v_x B_z$$

with $\tau' = \frac{\tau}{1+i\omega\tau}$

$$\Rightarrow v_x = \frac{q\tau'}{m^*} (E_x - \frac{q\tau'}{m^*} v_x B_z^2)$$

$$= \frac{q\tau'}{m^*} E_x - \frac{q^2 \tau'^2}{m^{*2}} v_x B_z^2$$

$$\Rightarrow v_x \left(1 + \frac{q^2 \tau'^2}{m^{*2}} B_z^2 \right) = \frac{q\tau'}{m^*} E_x$$

$$\Rightarrow v_x (1 + \omega_c^2 \tau'^2) = \frac{q\tau'}{m^*} E_x \quad \text{with } \omega_c = \frac{qB}{m^*}$$

$$\sigma = \frac{j_x}{E_x} = \frac{Nq v_x}{E_x} = Nq \frac{q\tau'}{m^*} \left(\frac{1}{1 + \omega_c^2 \tau'^2} \right)$$

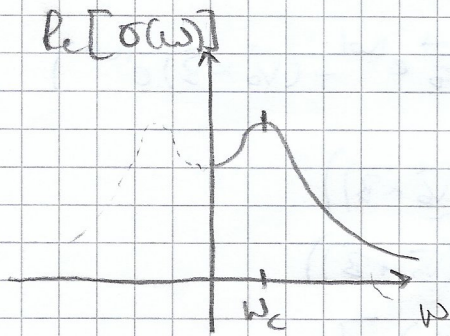
$$= \frac{Nq^2}{m^*} \frac{\tau}{1+i\omega\tau} \left(\frac{1}{1 + \omega_c^2 \frac{\tau^2}{(1+i\omega\tau)^2}} \right) = \frac{Nq^2 \tau}{m^*} \left(\frac{1}{(1+i\omega\tau) + \omega_c^2 \frac{\tau^2}{(1+i\omega\tau)}} \right)$$

$$= \frac{Nq^2 \tau}{m^*} \frac{1+i\omega\tau}{(1+i\omega\tau)^2 + \omega_c^2 \tau^2} = \frac{Nq^2 \tau}{m^*} \frac{1+i\omega\tau}{1 + (\omega_c^2 - \omega^2) \tau^2 + 2i\omega\tau}$$

Why neglect influence magn. field of Q-wave Super wave

$B \approx B_0$?
in this case given!

$$Re(\sigma) = \frac{1}{2}(\sigma + \sigma^*) = \frac{u g^2 \tau}{2u^*} \frac{[2(\omega_c^2 \tau^2 + \omega^2 \tau^2 + 1)]}{\sqrt{\omega_c^4 \tau^4 - 2\omega_c^2 \tau^6 \omega^2 + 2\omega_c^2 \tau^2 + \tau^4 \omega^4 + 2\tau^2 \omega^2 + 1}}$$



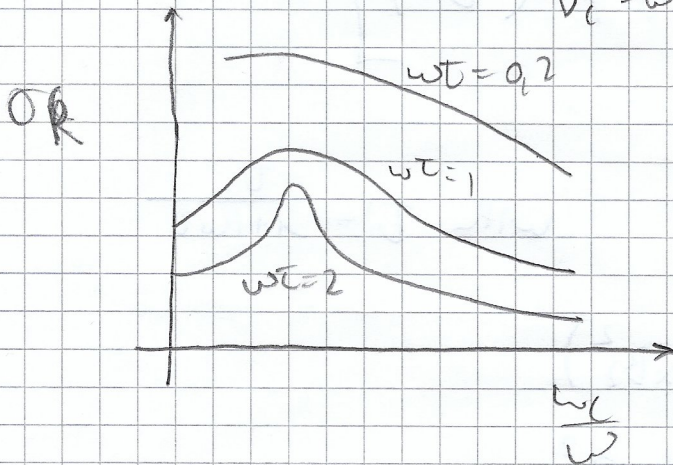
Nicer:

$$\frac{\sigma_R}{\sigma_0} = \frac{1 + v^2 + v_c^2}{(1 + v_c^2 - v^2)^2 + 4v^2}$$

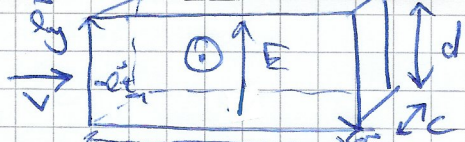
$$v = \omega \tau$$

$$\sigma_0 = \frac{Ne^2 \tau}{m^*}$$

$$v_c = \omega_c \tau$$



3. Hall effect



Instead
no force
Rotation so
that fields
in one direction
always allowed?

a) $F_L = q\vec{E} + q\vec{v} \times \vec{B} = 0$ with $\vec{B} = B\vec{e}_z$, $\vec{E} = E\vec{e}_y$, $\vec{v} = v\vec{e}_x$

$$\Rightarrow qE\vec{e}_y = -qvB(\vec{e}_x \times \vec{e}_z) = -qvB(-\vec{e}_y)$$

$$\Rightarrow E = vB$$

$$\Rightarrow \frac{U}{d} = vB$$

$$\Rightarrow U = e \cdot d \cdot vB$$

$$\Leftrightarrow U_H = \frac{1}{ne} IB$$

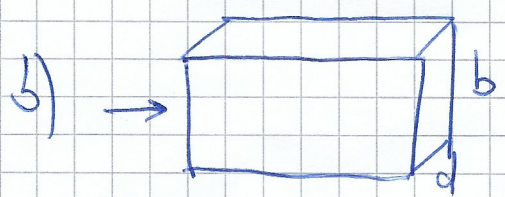
$$\Rightarrow U_H = \frac{1}{nec} B$$

$$|E = \frac{U}{d}$$

$$|d, c$$

$$|A = cd, \vec{I} = neAV$$

Cross-sectional area



$$d = 0,2 \text{ cm} \quad b = 1 \text{ cm}$$

$$n = 10^{15} \frac{1}{\text{cm}^3}$$

$$B = 100 \text{ mT}$$

$$U_H = 10 \text{ mV}$$

Volt meter
sensitivity
 $\approx 10 \mu\text{V}$?
need at least
this voltage

$$I = \frac{nec}{B} U_H = 32,04 \text{ A}$$

c) $R_H = \frac{1}{ne} \frac{\mu_n - \mu_p}{\mu_n + \mu_p}$, $U_H = R_H \frac{IB}{d}$

high temperatures, $n \approx n_i \gg N_D$ (intrinsic regime)

N_C, N_V : eff. density of e, holes

$$n_i = p_i = \sqrt{N_C N_V} \exp\left(-\frac{E_0}{2k_B T}\right)$$

intrinsic carrier concentration

$$R_H = \frac{1}{2e} \left(\frac{k_B T}{2\pi \hbar^2}\right)^{-3/2} (m_e m_h)^{-3/4} \exp\left(-\frac{E_0}{2k_B T}\right) \cdot \mu'$$

Plot
 $(R_H) T^{3/2}$ vs $\frac{1}{T}$
by $(R_H) T^{3/2} = \text{const} + \frac{E_0}{2k_B T}$

4. Properties of doped Silicon

$$N_D = \frac{10^{16}}{\text{cm}^3}, \quad n_i = 1,45 \cdot 10^{10} \frac{1}{\text{cm}^3}, \quad N_C = 2,8 \cdot 10^{19} \frac{1}{\text{cm}^3}$$

$T = 300\text{K}$

$$n \approx 2N_D \left(1 + \sqrt{1 + 4 \frac{N_D}{N_C} e^{\frac{E_D}{k_B T}}} \right)^{-1}, \quad E_D = E_C - E_D$$

intermediate temperatures with $4 \frac{N_D}{N_C} e^{\frac{E_D}{k_B T}} \ll 1$: Saturated regime:

$$n \approx N_D = 10^{16} \frac{1}{\text{cm}^3}$$

$$E_F(T) = E_C - k_B T \ln \left(\frac{N_C(T)}{N_D} \right), \quad E_C = 3,4\text{eV} \quad (\text{chapter 9, p 2})$$

(Fig. 9.1.5)

$$\Rightarrow E_F = 3,40\text{eV} - 0,21\text{eV} = 3,19\text{eV}$$