

Disclaimer

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<https://www.physics-and-stuff.com/>

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General relativity F. Home exercise

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$$05.06.2018 \text{ (H10)} \quad G(r) = \begin{cases} G_0 = \text{const} & , r < r_0 \\ 0 & , r > r_0 \end{cases}$$

(F17)
feed
2

✓
Radius r_0
free input
parameter -
what else?

meaning?

→ If we had
 ϵ and p , we could
compute G

Can choose

JK arbitrary
S.t. $G_{\mu\nu}T^{\mu\nu}$
fulfilled

$$g_{00} = e^{\nu(r)}, \quad g_{11} = -e^{\lambda(r)}, \quad g_{22} = -r^2, \quad g_{33} = -r^2 e^{2\nu(r)}$$

$$G_0^0 = \frac{1}{r^2} - \left(\frac{1}{r^2} - \frac{\lambda'(r)}{r} \right) e^{-\lambda(r)}$$

$$G_1^1 = \frac{1}{r^2} - \left(\frac{1}{r^2} + \frac{\nu'(r)}{r} \right) e^{-\lambda(r)}$$

$$G_2^2 = G_3^3 = -\frac{1}{2} (\nu''(r) + \frac{\nu'^2(r)}{2} - \frac{\lambda'(r)\nu'(r)}{2} + \frac{\nu' - \lambda'}{r}) e^{-\lambda(r)}$$

a)

$$G_0^0 = 8\pi G T_0^0 \Leftrightarrow 8\pi G \epsilon_0 = \frac{1}{r^2} - \left(\frac{1}{r^2} - \frac{\lambda'(r)}{r} \right) e^{-\lambda(r)}$$

$$\rightarrow \frac{dp}{dr} = (\epsilon + p) \frac{\partial}{\partial r} \left(\frac{1}{r^2} - \frac{\lambda'(r)}{r} \right) e^{-\lambda(r)}$$

$$\text{grav. inv. } \frac{dp}{dr} = 1 ?$$

$$\text{and } T^0_0 = \frac{1}{g_{00}} (\epsilon + p)$$

$$\rightarrow p g^0_0$$

$$\text{also if } g \text{ not diagonal?}$$

$$\Leftrightarrow 8\pi G \int_0^r \frac{dp}{dr} r^2 \epsilon_0 = 1 - (1 - \lambda'(r)) e^{-\lambda(r)} = 1 - e^{-\lambda(r)} + \lambda'(r) e^{-\lambda(r)}$$

$$\Leftrightarrow 8\pi G \int_0^r r^2 \epsilon_0 = r - r e^{-\lambda(r)} = r (1 - e^{-\lambda(r)})$$

$$= \frac{M(r)}{4\pi}$$

$$\Leftrightarrow 2GM(r) = r (1 - e^{-\lambda(r)})$$

$$\Leftrightarrow e^{-\lambda(r)} = 1 - \frac{2GM(r)}{r}$$

$$\Leftrightarrow e^{-\lambda(r)} = \left(1 - \frac{2GM(r)}{r} \right)^{-1}$$

$$r < r_0: M(r) = 4\pi \int_0^r r^2 \epsilon_0 = \frac{4\pi}{3} r^3 \epsilon_0$$

$$\rightarrow e^{-\lambda(r)} = \left(1 - \frac{8\pi G r^3 \epsilon_0}{3r} \right)^{-1} = \left(1 - \frac{r^2}{R^2} \right)^{-1}$$

Why not evaluate

M_0 as well?

→ can be

easily cal.

$$\rightarrow 1 - \frac{1}{2} \left(\frac{r_0^3}{R^2} \right)$$

$$R^2 = \frac{3}{8\pi G \epsilon_0}$$

$$r > r_0: M(r) = 4\pi \int_0^r r^2 \epsilon_0 = M_0$$

$$= \frac{4\pi r_0^3 \epsilon_0}{3}$$

$$\rightarrow e^{-\lambda(r)} = \left(1 - \frac{2GM_0}{r} \right)^{-1}$$

Q) Also had a third eq. arising from $T_{\mu\nu}^{\text{ext}}$ or the Einstein eq. :

(1)

$$-2p'(r) = v'(r)(p(r) + e(r))$$

$$\Rightarrow p'(r) = -\frac{v'(r)}{2}(p(r) + e_0) = -\frac{v'(r)}{2}p(r) - \frac{v'(r)}{2}e_0 \quad (*)$$

\hookrightarrow first solve hom. diff. eq. $p'(r) = -\frac{v'(r)}{2}p(r)$

in fact:
 $v' = \frac{-2p'}{e_0 + p}$
 $= -2 \log(e_0 + p)$

$\hookrightarrow p(r) = C e^{-\frac{v}{2}}$

within we the special solution $p(r) = -e_0$ for the inhom. diff eq. and add it to $p(r)$.

$$\hookrightarrow p(r) = C e^{-\frac{v}{2}} - e_0 \text{ solves } (*)$$

$$\Rightarrow p(r) + e_0 = C e^{-\frac{v}{2}}$$

Q) From $G_{00}^0 = 8\pi G_0 T_0^0$ and $G_{11}^1 = 8\pi G_1 T_1^1$, we find

$$8\pi G_0 = \frac{1}{r^2} - \left(\frac{1}{r^2} - \frac{\lambda''}{r} \right) e^{-\lambda}$$

$$-8\pi G_1 = \frac{1}{r^2} - \left(\frac{1}{r^2} + \frac{v'}{r} \right) e^{-\lambda}$$

$$\hookrightarrow 8\pi G_1(p + e_0) = \frac{\lambda' + v'}{r} e^{-\lambda} = 8\pi G_1 c e^{-\frac{v}{2}}$$

(2)

$$\hookrightarrow (\lambda' + v') e^{\frac{v}{2}} = 8\pi G_1 c r e^\lambda$$

$$\hookrightarrow v' e^{\frac{v}{2}} = 8\pi G_1 c r e^\lambda - \lambda' e^{\frac{v}{2}}$$

$$\begin{aligned} \hookrightarrow 2y' &= 8\pi G_1 c r e^\lambda - \lambda' y \\ y &= e^{\frac{v}{2}} \end{aligned}$$

use the ansatz $y = e^{\frac{v}{2}} = A - B(1 - \frac{r^2}{R^2})^{\frac{v}{2}}$ $\hookrightarrow y' = \frac{B\lambda'}{2} e^{-\lambda/2}$

$$= A - B e^{-\lambda/2}$$

$$\hookrightarrow B\lambda' e^{-\lambda/2} = 8\pi G_1 c r e^\lambda - A\lambda' + B\lambda' e^{-\lambda/2}$$

$$\hookrightarrow 8\pi G_1 c r e^\lambda = A\lambda' \Leftrightarrow c = \frac{A\lambda'}{8\pi G_1 r} e^{-\lambda}$$

$$\left| e^{-\lambda} = \left(1 - \frac{r^2}{R^2}\right)^{\frac{v}{2}} \right. \Leftrightarrow -\lambda' e^{-\lambda} = -\frac{2r}{R^2} \Leftrightarrow \lambda' e^{-\lambda} = \frac{2r}{R^2}$$

$$\hookrightarrow c = \frac{2A}{8\pi G_1 R^2}$$

d)

$$\text{From b)} \quad p(r) = C e^{-\frac{r}{R}} - E_0$$

$$\text{From a)} \quad R^2 = \frac{3}{8\pi G E_0} \Leftrightarrow E_0 = \frac{3}{8\pi G R^2}$$

(1)

Using c), we find

$$\begin{aligned} p(r) &= C e^{-\frac{r}{R}} - E_0 = \frac{C}{A - 3B\sqrt{1 - \frac{r^2}{R^2}}} - \frac{3}{8\pi G R^2} \\ &= \frac{8\pi G R^2 C - 3A + 3B\sqrt{1 - \frac{r^2}{R^2}}}{8\pi G R^2 \left\{ A - B\sqrt{1 - \frac{r^2}{R^2}} \right\}} \\ \text{from c)} &= \frac{3B\sqrt{1 - \frac{r^2}{R^2}} - A}{A - B\sqrt{1 - \frac{r^2}{R^2}}} \end{aligned}$$

(1)

e) To fix A and B, we use $p(r_0) = 0$ and
 $e^{V(r_0)} = 1 - \frac{2Gm_0}{r_0}$

i.e. continuity of p and e^V at the boundary of the star, as from the lecture

$$e^{V(r)} = \left(1 - \frac{2Gm(r)}{r}\right) \text{ for } r > r_0$$

$$p(r) = 0 \text{ for } r > r_0$$

$$\Rightarrow p(r_0) = \frac{1}{8\pi G R^2} \frac{3B\sqrt{1 - \frac{r_0^2}{R^2}} - A}{A - B\sqrt{1 - \frac{r_0^2}{R^2}}} = 0$$

$$\Leftrightarrow A = 3B\sqrt{1 - \frac{r_0^2}{R^2}} = 3Be^{-\lambda(r_0)/2}$$

$$e^{\lambda(r_0)/2} = A - B\sqrt{1 - \frac{r_0^2}{R^2}} \stackrel{!}{=} \sqrt{1 - \frac{2Gm_0}{r_0}}$$

$$= A - \frac{A}{3} = \frac{2A}{3}$$

$$\Leftrightarrow A = \frac{3}{2}\sqrt{1 - \frac{2Gm_0}{r_0}}$$

$$\Rightarrow B = \frac{A}{3\sqrt{1 - \frac{r_0^2}{R^2}}} = \frac{\sqrt{1 - \frac{2Gm_0}{r_0}}}{2\sqrt{1 - \frac{r_0^2}{R^2}}}$$

$p(r_0) \geq 0$ as well on the boundary itself?

We don't know exactly what it's at the boundary microscopically, but macroscopically take $p(r_0 + \delta)$ for $\delta \rightarrow 0$, which can be assumed to vanish.

Can those A, B be simplified further?
 Who did that in tutorial?